

Towards three-loop QCD corrections to the time-like splitting functions

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in collaboration with

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Three-loop time-like $q \rightarrow g$ splitting function



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Why three-loop?

Because since 1980 two-loop splitting functions were calculated by various methods probably as more times as any other expression.

Space-like

- ▶ axial gauge (PV prescription) [Curci, Furmanski, Petronzio '80](#); [Ellis, Vogelsang '98](#)
- ▶ Feynman gauge (x -space) [Floratos, Kounnas, Lacaze '81](#)
- ▶ axial gauge (ML prescription) [Bassetto, Heinrich, Kunszt, Vogelsang '98](#)
- ▶ Feynman gauge (Mellin space) [Moch, Vermaseren '99](#)
- ▶ axial gauge (NPV prescription) [OG, Jadach, Skrzypek, Kusina '14](#) ([A.Kusina's talk](#))

Time-like

- ▶ axial gauge (PV prescription) [Furmanski, Petronzio '80](#)
- ▶ Feynman gauge (x -space) [Floratos, Kounnas, Lacaze '81](#)
- ▶ **analytic continuation** [Stratmann, Vogelsang '96](#);
[Blumlein, Ravindran, van Neerven '00](#); [Moch, Vogt '07](#)
- ▶ Feynman gauge (Mellin space) [Mitov, Moch '06](#)

Three-loop time-like $q \rightarrow g$ splitting function

Why time-like?

Because space-like three-loop splitting functions are already calculated. Moch, Vermaseren, Vogt '04

(*Nuclear Physics B* selected top paper in 2004)

But can't one use some trick to derive them from space-like results?

Examples of tricks (analytic continuation): Drell, Levy, Yan '70; Gribov, Lipatov '72

Sure!

- ▶ NNLO non-singlet Mitov, Moch, Vogt '06
- ▶ NNLO singlet ($q \rightarrow q$ and $g \rightarrow g$) Moch, Vogt '07
- ▶ NNLO singlet ($q \rightarrow g$ and $g \rightarrow q$) Almasy, Moch, Vogt '11

Three-loop time-like $q \rightarrow g$ splitting function should be calculated directly.

1. Mass factorization (algebraic) relations

The unpolarized differential cross-section for " $\gamma^* \rightarrow$ partons" decay

$$\frac{1}{\sigma_{\text{tot}}} \frac{d^2\sigma^H}{dx d\cos\theta} = \frac{3}{8}(1 + \cos^2\theta) \mathcal{F}_T(x) + \frac{3}{4}\sin^2\theta \mathcal{F}_L(x) + \frac{3}{4}\cos\theta \mathcal{F}_A(x)$$

The mass factorization relations are [Vermaseren, Vogt, Moch '05](#)

$$\mathcal{F}_T^{(1)} = -\frac{2}{\epsilon} P_{gq}^{(0)} + c_{T,g}^{(1)} + \epsilon a_{T,g}^{(1)} + \epsilon^2 b_{T,g}^{(1)}$$

$$\begin{aligned} \mathcal{F}_T^{(2)} = & \frac{1}{\epsilon^2} \left\{ P_{gq}^{(0)} \left(P_{qq}^{(0)} + P_{gg}^{(0)} + \beta_0 \right) \right\} + \frac{1}{\epsilon} \left\{ P_{gq}^{(1)} + 2c_{T,q}^{(1)} P_{gq}^{(0)} + c_{T,g}^{(1)} P_{gg}^{(0)} \right\} \\ & + c_{T,g}^{(2)} - 2a_{T,q}^{(1)} P_{gq}^{(0)} - a_{T,g}^{(1)} P_{gg}^{(0)} + \epsilon \left\{ a_{T,g}^{(2)} - 2b_{T,q}^{(1)} P_{gq}^{(0)} - b_{T,g}^{(1)} P_{gg}^{(0)} \right\} \end{aligned}$$

$$\begin{aligned} \mathcal{F}_T^{(3)} = & -\frac{1}{6\epsilon^3} \left\{ P_{gi}^{(0)} P_{ij}^{(0)} P_{jg}^{(0)} + 3\beta_0 P_{gi}^{(0)} P_{ig}^{(0)} + 2\beta_0^2 P_{gg}^{(0)} \right\} \\ & + \frac{1}{6\epsilon^2} \left\{ 2P_{gi}^{(0)} P_{ig}^{(1)} + P_{gi}^{(1)} P_{ig}^{(0)} + 2\beta_0 P_{gg}^{(1)} + 2\beta_1 P_{gg}^{(0)} + 3P_{gi}^{(0)} \left(P_{ij}^{(0)} + \beta_0 \delta_{ij} \right) c_{\phi,j}^{(1)} \right\} \\ & - \frac{1}{6\epsilon} \left\{ 2P_{gg}^{(2)} + 3P_{gi}^{(1)} c_{\phi,i}^{(1)} + 6P_{gi}^{(0)} c_{\phi,i}^{(2)} - 3P_{gi}^{(0)} \left(P_{ij}^{(0)} + \beta_0 \delta_{ij} \right) a_{\phi,j}^{(1)} \right\} \\ & + c_{\phi,g}^{(3)} - \frac{1}{2} P_{gi}^{(1)} a_{\phi,i}^{(1)} - P_{gi}^{(0)} a_{\phi,i}^{(2)} + \frac{1}{2} P_{gi}^{(0)} \left(P_{ij}^{(0)} + \beta_0 \delta_{ij} \right) b_{\phi,j}^{(1)} + O(\epsilon) \end{aligned}$$

2. Fragmentation Functions

$$\mathcal{F}_T(x, \epsilon) = \frac{2}{2-d} \left(\frac{k_0 \cdot q}{q^2} W_\mu^\mu + \frac{k_0^\mu k_0^\nu}{k_0 \cdot q} W_{\mu\nu} \right)$$

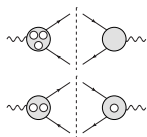
- ▶ The hadronic tensor is defined as

$$W_{\mu\nu}(x, \epsilon) = \frac{x^{d-3}}{4\pi} \int d\text{PS}(n) M_\mu(n) M_\nu(n)$$

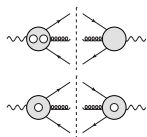
- ▶ $d\text{PS}(n)$ is n -particle real phase-space
- ▶ amplitude $M^\mu(n)$ describes process " $\gamma^*(q) \rightarrow g(k_0) + n$ partons"

$$\begin{aligned} \mathcal{F}_T^{(1)} &= \text{diagram 1} + \text{diagram 2} \\ \mathcal{F}_T^{(2)} &= \text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \text{diagram 4} \\ \mathcal{F}_T^{(3)} &= \text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \text{diagram 4} + \text{diagram 5} + \text{diagram 6} + \text{diagram 7} + \text{diagram 8} \end{aligned}$$

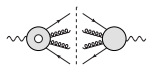
3. Final Considerations for NNLO corrections



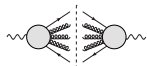
Pure-virtual contributions contain overall $\delta(1-x)$ factor.
We do not consider such contributions.



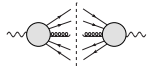
Can be extracted from [Garland, Gehrmann, Glover, Koukoutsakis, Remiddi '01](#)
Calculated by [Duhr, Gehrman, Jaquier 1411.3587 \[hep-ph\]](#)



One-loop helicity amplitudes by [Bern, Dixon, Kosower '97](#)
Final-state integration is of NLO complexity — simple.

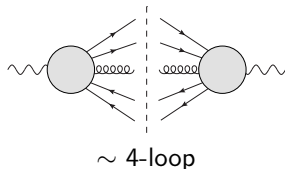
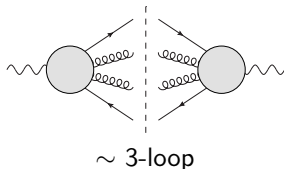


Contribution is known from analytical continuation by [Almasy, Moch, Vogt '11](#)



Unknown!

Final-state integration



The main challenge of the calculation is n -particle final-state integration:

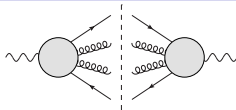
$$\int d\text{PS}(n) = \int \prod_{i=0}^n d^m k_i \delta^+(k_i^2) \delta\left(x - \frac{2k_0 \cdot q}{q^2}\right) \delta\left(q - \sum_{j=0}^n k_j\right)$$

**Since recently such integrals can be found
with Integration-By-Parts method.**

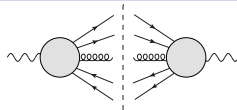
Software we consider:

- ▶ LiteRed [Lee '13](#)
- ▶ Reduze [von Manteuffel, Studerus '12](#)

Preparation



▶ 8 amplitudes



▶ 48 amplitudes

QGraf

FORM

- ▶ trace of gamma matrices
- ▶ index contraction
- ▶ color traces
- ▶ partial fractioning

Mathematica

- ▶ analyze symmetries
- ▶ split by topologies

LiteRed

- ▶ find IBP reduction rules
- ▶ find masters

▶ 499 integrals

▶ ~10 h

▶ 9 masters

▶ 55 614 integrals

▶ ~1000 h (multi-thread)

▶ ~80 masters

System of Differential Equations for Masters at NLO

$$\left(\begin{array}{cccccccccccc} \frac{(2\epsilon-1)(2x-1)}{(1-x)x} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \frac{3\epsilon-2}{(1-x)x} & -\frac{3\epsilon-1}{x} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \frac{(2\epsilon-1)(3\epsilon-1)}{\epsilon(x-1)x} & \frac{2\epsilon}{1-x} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \frac{(2\epsilon-1)(3\epsilon-2)(3\epsilon-1)}{2\epsilon(x-1)x^2} & \frac{(2\epsilon-1)(3\epsilon-1)(x^2-10x+1)}{2(1-x)x^2(x+1)} & 0 & \frac{2\epsilon(x^2-3x-2)}{(1-x)x(x+1)} & \frac{2\epsilon(6\epsilon-1)}{(1-x)x} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & \frac{(2\epsilon-1)(3\epsilon-1)}{\epsilon(x-1)x} & 0 & \frac{2}{x-1} & \frac{6\epsilon-1}{1-x} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ \frac{4(2\epsilon-1)(3\epsilon-2)(3\epsilon-1)}{\epsilon^2(1-x)x^3(x+1)} & \frac{4(2\epsilon-1)(3\epsilon-1)(x^2-x+1)}{\epsilon(x-1)x^3(x+1)^2} & 0 & \frac{4(x^2+1)}{(x-1)x^2(x+1)^2} & \frac{2(6\epsilon-1)}{(1-x)x^2(x+1)} & \frac{(2\epsilon+1)(2x+1)}{-x(x+1)} & 0 & 0 & 0 & 0 & 0 & 0 \\ \frac{2(2\epsilon-1)(3\epsilon-2)(3\epsilon-1)}{\epsilon^2(x-1)^2x^2} & -\frac{(2\epsilon-1)(3\epsilon-1)}{\epsilon(x-1)x^2} & \frac{2\epsilon}{(x-1)x} & 0 & 0 & 0 & 0 & \frac{4\epsilon+1}{-x} & 0 & 0 & 0 & 0 \\ \frac{2(1-2\epsilon)(3\epsilon-2)(3\epsilon-1)}{\epsilon^2(x-1)^2x^3} & \frac{2(2\epsilon-1)(3\epsilon-1)(3x-1)}{\epsilon(x-1)^3x^3} & \frac{4\epsilon}{(1-x)^2x} & \frac{4(x^2+1)}{(x-1)^3x^2} & \frac{2(6\epsilon-1)(x+1)}{(1-x)^3x^2} & 0 & 0 & 0 & \frac{(2\epsilon+1)(2x-1)}{(1-x)x} & 0 & 0 & \frac{(2\epsilon+1)(2x-1)}{(1-x)x} \end{array} \right)$$

Following properties ensure that solution can be found in terms of **Harmonic Polylogarithms** Remiddi, Vermaseren '99:

- ▶ Alphabet (letters): $\{x, 1-x, 1+x\}$
- ▶ Main diagonal contains only $\left\{\frac{1}{x}, \frac{1}{1-x}, \frac{1}{1+x}\right\}$ terms
- ▶ **Non-triangular terms vanish** in $\epsilon \rightarrow 0$ limit

Solutions for Masters at NLO

Masters are found in terms of HPLs
as ϵ -series to any order.

$$M_2(x) = M_2^{(0)} + \epsilon M_2^{(1)} + \epsilon^2 M_2^{(2)} + \dots$$

$$M_2^{(0)} = x^{1-3\epsilon} (C_2^{(0)} - 2 \ln x C_1^{(0)})$$

$$M_2^{(1)} = x^{1-3\epsilon} (C_2^{(1)} - 2 \ln x C_1^{(1)} + (3H_0 - 4H_2 - 2H_{0,0}) C_1^{(0)})$$

$$M_2^{(2)} = x^{1-3\epsilon} (C_2^{(2)} - 2 \ln x C_1^{(2)} + (3H_0 - 4H_2 - 2H_{0,0}) C_1^{(1)} \\ + (6H_2 - 4H_3 + 3H_{0,0} - 4H_{2,0} - 8H_{2,1} - 2H_{0,0,0}) C_1^{(0)})$$

Actual challenge is to find integration constants $C_n^{(k)}$!

To be continued...

Summary

Done:

- ▶ **Preparation steps** to start integration
→ code in QGRAF, FORM, Mathematica
- ▶ **Simplify** integrals
→ various symmetries and partial fractioning
- ▶ **IBP identities** for NNLO case
→ reduces 55 614 integrals to just ~ 80 masters
→ $\sim 1000\text{h}$ and still running
(CC1 cluster at IFJ PAN & Phenod cluster at DESY)
- ▶ **Solution for masters**
→ from differential equations to any order in ϵ -parameter

In progress:

- ▶ Implementation of the **boundary conditions finder**
- ▶ Three-loop time-like $q \rightarrow g$ splitting function